ment, within experimental errors, with the values observed in the s-facial isomer (Kobayashi, Marumo & Saito, 1972). All the observed Co-N distances are normal for Co(III) complexes of linear aliphatic polyamines. The chelate rings formed by a dien ligand in the complex cation A have eclipsed envelope and symmetrical skew conformations, while those in the complex cation B are both eclipsed envelope conformations. They are shown in Fig. 3. The Newman projections along the C-C bonds are illustrated in Fig. 4. All the dihedral angles N-C-C-N are in the range $43 \sim 45^{\circ}$ and are much smaller than that in a typical gauche structure. The bond distances and angles in the hexacyanocobaltate(III) anion are listed in Table 5. They may be compared with those observed in $(-)_{589}$ -[Co(penten)] [Co(CN)₆] 2H₂O (Muto, Marumo & Saito, 1970). The closest approach between complex cations and anions occurs between the terminal nitrogen atom of the cation and the nitrogen atom of the anion: $N(3) \cdots N(12)$ and $N(4) \cdots N(9)$ are 2.918 and 3.086 Å respectively. The positions of hydrogen atoms attached to N(3) and N(4) suggest that N(3) \cdots N(12) and N(4) \cdots N(9) are hydrogen bonds. The water molecules are surrounded by two nitrogen atoms of the anion and a terminal nitrogen atom of the cation as indicated by broken lines in Fig. 2. These $N \cdots O$ distances are mostly less than 3 Å. Other interatomic distances less than 3.5 Å between complexes are listed in Table 6.

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Fig. 4. Newman projections along the C-C bonds in the chelate ring.

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The Crystal Structure of CsMnCl₃

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The crystal structure of CsMnCl₃, determined from Weissenberg data, is trigonal, space group R3m. The hexagonal cell contains 9 molecules and has dimensions a = 7.29 (1) and c = 27.48 (5) Å. The c parameter accommodates 9 close-packed layers of composition CsCl₃ and the Mn ions are octahedrally coordinated by Cl ions. The structure is closely related to that of CsNiCl₃.

Introduction

Many complex halides of composition CsMCl₃, where M represents a metal ion in the first transition series, have the same structure as CsNiCl₃ (Tishchenko, 1955), which consists of hexagonally close-packed CsCl₃ layers with Ni octahedrally coordinated by Cl ions such that the Ni–Cl octahedra share opposite faces to form infinite chains of composition $(\text{NiCl}_3)_n^n$ parallel to **c**, the space group being $P6_3/mmc$ with $a \simeq 7.2$ Å, $c/a \simeq 0.82$ and Z = 2. Recent work (Goodyear & Kennedy, 1972) has shown that CsMnBr₃ also has this structure.

An investigation of the unit cell of CsMnCl₃ by Kes-

tigian, Croft & Leipziger (1967) using a powder diffractometer technique indicated a hexagonal cell with a =7.288 and c = 27.44 Å, and an observed density of 3.48 g cm⁻³ required 9 molecules per unit cell. The *a* parameter of this cell is very similar to that of CsNiCl₃ and suggests that this structure also contains CsCl₃ closepacked layers; but the axial ratio, which is approximately $4\frac{1}{2}$ times that of the CsNiCl₃ cell, would require a different stacking sequence of such layers.

In the present work, a Weissenberg study of $CsMnCl_3$ has confirmed the unit-cell data found by Kestigian, Croft & Leipziger and a consideration of a suitable stacking arrangement of $CsCl_3$ layers to explain the large *c* parameter has enabled the crystal structure to be determined.



Fig. 1. Close-packed CsCl₃ layers classified according to the position of the unit cell (broken line). The circles and squares represent Cl and Cs atoms respectively. The triangles show possible sites for Mn atoms between different pairs of layers; ab, bc and ac are sites between types A and B, B and C, A and C respectively.

Experimental

Single crystals of CsMnCl₃ were prepared by heating an equimolar mixture of MnCl₂ and CsCl in a sealed evacuated silica tube to a temperature above the melting point and then cooling the specimen to room temperature at a rate of about 5°C per hour. Because of the extremely hygroscopic nature of the material, single crystals suitable for X-ray examination were selected and examined in a stream of dry nitrogen and were then mounted in Lindemann-glass capillary tubes containing P_2O_5 .

The unit-cell dimensions were determined from Cu $K\alpha$ Weissenberg photographs, taken about the *a* axis, using the method of Main & Woolfson (1963). The density was determined by weighing a sample of the material quickly in air and in toluene, and agreement between the observed and calculated densities was obtained by assigning 9 molecules of CsMnCl₃ per unit cell. The complete crystal data are shown in Table 1.

Table 1. Crystal data

Formula, CsMnCl ₃	F.W. 294·20
Hexagonal cell	a = 7.29 (1), $c = 27.48$ (5) Å
-	$Z=9, D_0=3.48, D_x=3.48 \text{ g cm}^{-3}$
Mean r (cm)	0.0037
μ (cm ⁻¹)	$102 \ (\lambda = 0.7107 \ \text{\AA})$

Intensity data were collected from equi-inclination Weissenberg photographs taken about the *a* axis with Mo $K\alpha$ radiation. The intensities of 687 reflexions from layer lines 0 to 6 were measured from multiple film exposures using a Joyce-Loebl flying-spot microdensitometer, although the intensities of a few very weak reflexions had to be estimated visually against a calibrated scale; 251 reflexions were unobserved. These data were corrected for the Lorentz-polarization factor and for absorption using the factors given by Bond (1959) for a cylindrical specimen.

The observed reflexions all satisfied the condition -h+k+1=3n, pointing to a rhombohedral primitive cell and $R\overline{3}m$, R3m, R32, $R\overline{3}$ and R3 as possible space groups.

Determination and description of the structure

The initial structure was determined by considering a stacking arrangement of $CsCl_3$ layers which would explain a c parameter of the order of 27 Å. In the $CsNiCl_3$ structure, the c parameter accommodates two close-packed layers, indicating a layer separation of about 3 Å and thus suggesting 9 layers in the unit cell of $CsMnCl_3$. This is also consistent with the assignment of 9 molecules per cell.

Since each Cs atom must be in contact with 12 Cl atoms, there are three possible positions of a layer in relation to the outline of the unit cell; these correspond

to the types A, B and C of Fig. 1. The stacking sequence for a 9 layer structure would then be

 $BABACACBC \cdots$, or $(chh)_3$ in the usual nomenclature. In the CsNiCl₃ structure the layers are hexagonally close-packed, *i.e.* $BABA \cdots$.

In a close-packed array of $CsCl_3$ layers only one quarter of the octahedral sites are bounded exclusively by Cl atoms and each of these must be occupied by a Mn atom in order to obtain the correct compositional formula. The triangular motifs in Fig. 1 indicate these sites in the unit cell between different pairs of layers; for example that labelled *ab* is such a site between *A* and *B*-type layers.

The idealized structure based on this reasoning is shown in Fig. 2, which is drawn in three parts to show the relation between the proposed structure and that of CsNiCl₃. The section between $z = -\frac{1}{6}$ and $z = \frac{1}{6}$ consists of a slice of CsNiCl₃-type structure $1\frac{1}{2}$ unit cells high. The next section, between $z = \frac{1}{6}$ and $z = \frac{1}{2}$, is a similar slice which is displaced -a/3 parallel to X and a/3 parallel to Y, whilst the section between $z = \frac{1}{2}$ and $z = \frac{5}{6}$ represents a third similar slice displaced a/3 parallel to X and -a/3 parallel to Y relative to the first slice.

The idealized positional parameters are consistent with the most symmetrical of the possible space groups, $R\overline{3}m$, and initially the model was tested by means of a block-diagonal three-dimensional least-squares refinement in this space group using individual isotropic temperature factors with Cs at equipoints 6(c) and 3(b), Mn at 6(c) and 3(a) and Cl at 18(h) and 9(d). In these calculations the atomic scattering factors given in *International Tables for X-ray Crystallography* (1962) for Cs⁺, Mn²⁺ and Cl⁻ were used and interlayer scaling was achieved by scaling the observed to the calculated structure factors. With the idealized positional coordinates and temperature factors similar to those found for the

Table 2. Refined temperature factors in the space group $R\overline{3}m$

	Equipoint	Coordinates	B (Å ²)
Mn(1)	3(a)	0,0,0	2.99
Mn(2)	6(c)	0, 0, z	0.38
Cs(1)	3(b)	$0, 0, \frac{1}{2}$	4.18
Cs(2)	6(c)	0, 0, z	1.04
C(1)	9(d)	1 ,0, 1	3.43
C(2)	18(h)	x.2x.z	0.62

CsMnBr₃ structure, the initial R value, $\sum ||F_o| - |F_c||/\sum |F_o|$, was as low as 20%. After about 10 cycles of refinement R decreased to 11.6%, at which stage







Г	abl	le	3.	Final	atomic	parameters
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Standard deviations are given in parentheses.

	Equipoint	x/a	<i>y</i> / <i>b</i>	z/c	<i>B</i> (Å ²)
Mn(1)	3(a)	0	0	-0.0002(7)	1.17 (8)
Mn(2)	3(a)	0	0	0.1160 (3)	1.34 (12)
Mn(3)	3(a)	0	0	-0.1159(3)	0.89 (9)
Cs(1)	3(<i>a</i>)	0	0	0.2817 (1)	2.09 (5)
Cs(2)	3(a)	0	0	-0.2816(1)	2.04 (5)
Cs(3)	3(a)	0	0	0.4968 (2)	1.86 (4)
Cl(1)	9(<i>b</i>)	0.1608 (6)	0.3216 (6)	0.0571 (3)	1.34 (9)
Cl(2)	9(b)	-0·1558 (5)	-0.3116(5)	-0·0573 (2)	0.97 (8)
Cl(3)	9(b)	-0.1606 (9)	-0·3212 (9)	0.1675 (5)	1.38 (7)

the shifts in the positional parameters were all less than $\frac{1}{10}$ of the estimated standard deviations. However, the refined temperature factors were not too satisfactory, since for each type of atom the temperature factor of an atom in a fixed position differed greatly from that of an atom in a position for which z could vary, as can be seen in Table 2.

Table 4. Magnitudes of observed and calculated structure factors

It was therefore decided to refine the structure in the non-centred space group R3m (No. 160) to allow shifts in the z parameter for each atom. With Cs and Mn each in three sets at (0,0,z) and Cl in three sets at (x, 2x, z), several cycles of refinement gave a minimum R value of 6.4%. The temperature terms were now much more reasonable, the shifts in all the parameters were less than $\frac{1}{3}$ of the estimated standard deviations and the calculated structure factors of each unobserved reflexion were less than the minimum observable value. Finally, an attempt was made at further refinement in the least symmetrical of the possible space groups, R3, to permit the x and y parameters of each chlorine atom to vary independently, but this did not change any of the parameters by more than one standard deviation. It was concluded that R3m was the correct space group. The final atomic parameters are given in Table 3 and the magnitudes of the calculated and observed structure factors compared in Table 4.

Bond lengths and angles are listed in Table 5. The average Mn–Cl and Cl–Cl distances in the Mn–Cl octahedra are 2.55 and 3.60 Å respectively, and the average Cs–Cl distance is 3.70 Å. These magnitudes and the fact that Pauling's electrostatic-valency rule is obeyed suggest that the structure is essentially ionic. The average Mn–Cl and Cl–Cl distances agree very well with those (2.56 and 3.61 Å respectively) found for the Mn–Cl octahedra in the structure of Na_2MnCl_4 by Goodyear, Ali & Steigmann (1971).

Each Mn(1)-Cl octahedron shares opposite faces with a Mn(2)-Cl and a Mn(3)-Cl octahedron and is thereby trigonally distorted owing to the mutual repulsion of the three neighbouring Mn^{2+} ions. The average separation of these ions, 3·18 Å, is significantly smaller than the corresponding separation, 3·26 Å, in the CsMnBr structure and this consequently leads to a greater degree of trigonal distortion than is found in the Mn-Br octahedron in the bromide; in fact, the distortion appears to be sufficiently great to prevent the formation of infinite chains of face-sharing octahedra.

Although Mn(2) and Mn(3) ions are equidistant from neighbouring Mn(1) ions within one standard deviation, the separation of planes (perpendicular to **c**) of neighbouring Mn(2) and Mn(3) ions is significantly shorter than the mean Mn(1)–Mn(2, 3) distance, 2·79 Å compared with 3·18 Å. Three corners of each Mn(3)– Cl octahedron are shared with Mn(2)–Cl octahedra, the Cl(3) ions being located at these common vertices. In the Mn(2)–Cl octahedron the Cl(3) ions are displaced towards the trigonal axis of the octahedron, and consequently away from the trigonal axis of a shared Mn(3)–Cl octahedron. The resulting Cl(3)–Cl(3) distances account for the difference in the distortion in these two types of octahedron.

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Table 5. Bond lengths and angles

IV	runnphen	y		
Mn(1)-Cl octahedron				
$Mn - Cl(1^i)$	3	2·570 (14) Å	$Cl(1^{i})-Mn-Cl(1^{i})$	86·3 (4)°
$Mn - Cl(2^i)$	3	2.516 (13)	$Cl(2^{i})-Mn-Cl(2^{i})$	85.3 (4)
$Cl(1^{i})-Cl(1^{i})$	3	3.517 (14)	$Cl(1^{i})-Mn-Cl(2^{i})$	94·2 (4)
$Cl(2^i)-Cl(2^i)$	3	3.408 (12)		
$Cl(1^{i})-Cl(2^{i})$	6	3.726 (11)		
Mn(2)-Cl octahedron				
$Mn - Cl(1^i)$	3	2.596 (10)	$Cl(1^{i})-Mn-Cl(1^{i})$	85.3 (2)
$Mn - Cl(3^i)$	3	2.472 (13)	$Cl(3^{i})-Mn-Cl(3^{i})$	90.5 (4)
$Cl(1^i)-Cl(1^i)$	3	3.517 (14)	$Cl(1^{i})-Mn-Cl(3^{i})$	92.1 (3)
$Cl(3^{1})-Cl(3^{1})$	3	3.512 (20)		
$Cl(1^{i})-Cl(3^{i})$	6	3.649 (16)		
Mn(3)-Cl octahedron	i			
$Mn - Cl(2^i)$	3	2.543 (8)	$Cl(2^i)-Mn-Cl(2^i)$	84.2 (2)
$Mn - Cl(3^{ii})$	3	2.577 (13)	$Cl(3^{ii})-Mn-Cl(3^{ii})$	94.3 (4)
$Cl(2^i) - Cl(2^i)$	3	3.408 (12)	$Cl(2^i) - Mn - Cl(3^{ii})$	90.6 (3)
$Cl(3^{ii})-Cl(3^{ii})$	3	3.778 (20)		
$Cl(2^{i}) - Cl(3^{ii})$	6	3.639 (15)		
Cs(1)-Cl distances				
$Cs-Cl(1^i)$	3	3.699 (9)		
$Cs-Cl(2^i)$	6	3.651 (5)		
$Cs-Cl(3^{ii})$	3	3.736 (14)		
Cs(2)-Cl distances				
$Cs-Cl(1^i)$	6	3.649 (5)		
$Cs-Cl(2^i)$	3	3.741 (8)		
Cs-Cl(3 ⁱ)	3	3.857 (14)		
Cs(3)-Cl distances				
$Cs-Cl(1^i)$	3	3.645 (10)		
$Cs-Cl(2^{11})$	3	3.820 (9)		
Cs-Cl(3 ⁱ)	6	3.648 (5)		

Standard deviations, attributable to e.s.d.'s in the positional and cell parameters, are given in parentheses.

Location of atoms [see Fig. 2(a) and (b)]:

Mn(1) at (0,0,0); Mn(2) at (0,0, $\frac{1}{9}$); Mn(3) at (0,0, $-\frac{1}{9}$). Cs(1) at $(\frac{1}{3}, \frac{2}{3}, -\frac{1}{18})$; Cs(2) at $(\frac{2}{3}, \frac{1}{3}, \frac{1}{18})$; Cs(3) at $(\frac{1}{3}, \frac{2}{3}, \frac{1}{6})$. Cl(1ⁱ) at $z = \frac{1}{18}$; Cl(2ⁱ) at $z = -\frac{1}{18}$; Cl(2ⁱⁱ) at $z = \frac{5}{18}$; Cl(3ⁱ) at $z = \frac{1}{6}$; Cl(3ⁱⁱ) at $z = -\frac{1}{6}$.

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